

6.2.1 Three-Dimensional Case

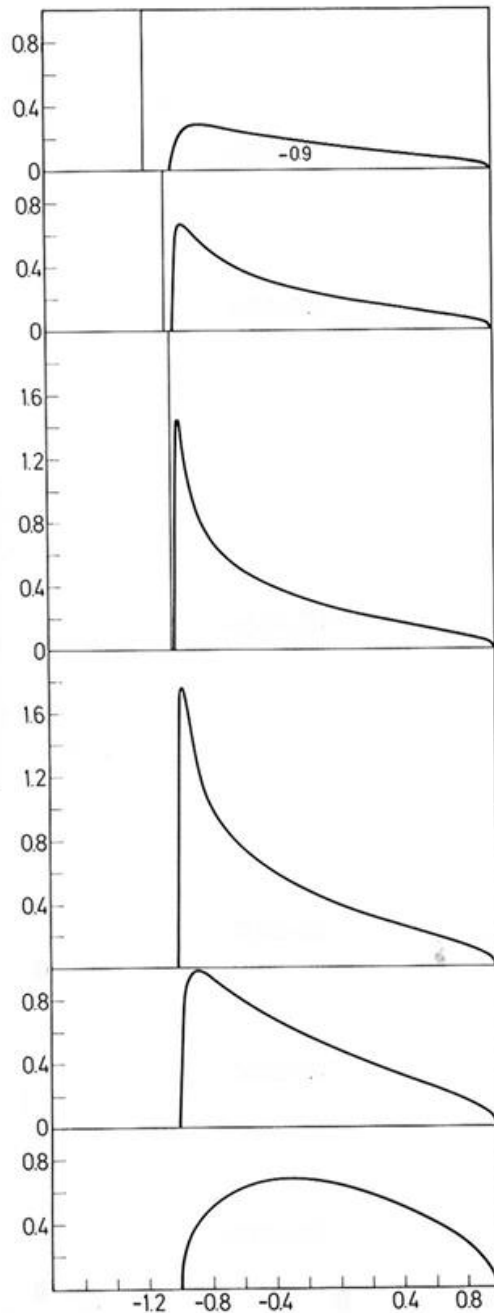


Fig.6.3. The perturbed DOS $\rho(\underline{\xi}; E)$ vs. E for different values of the attractive local potential. As $|\epsilon|$ increases from zero, states are pushed towards the lower band edge and at a critical value ($|\epsilon|/B = 0.5$) a discrete level is split off the continuum. The unperturbed DOS $\rho_0(\underline{\xi}; E)$ has been taken as in (5.55)

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Substituting in (6.9) we obtain

$$E_p - E_\ell \approx -C \exp\left(-\frac{1}{|\epsilon|\rho_d}\right) ; \quad \epsilon \leq 0 \quad . \quad (6.44)$$

Thus the binding energy $|E_p - E_\ell|$ for weak perturbation depends exponentially on the strength of the perturbation. This dependence stems directly from the discontinuity of the unperturbed DOS at the band edge, which in turn is a characteristic feature of two-dimensional systems. For the case of the

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Thus both I_ℓ and I_u are infinite, and consequently there is always a bound state no matter how small $|\epsilon|$ is. For small negative ϵ the bound level is very close to the band edge and thus one can use expression (6.47) in (6.9). We then obtain for the binding energy E_b the following expression

$$E_b = |E_p - E_\ell| \xrightarrow{\epsilon \rightarrow 0^-} \epsilon^2 \pi^2 C^2 \quad . \quad (6.48)$$

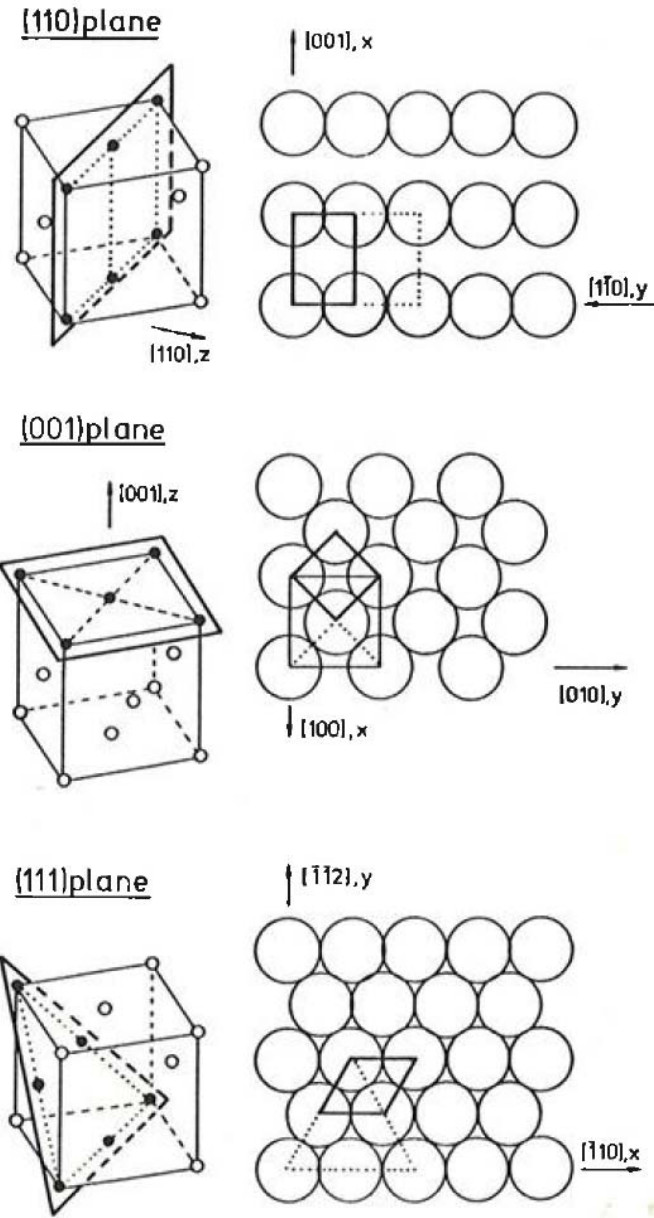


Fig.8.5. (100), (001) and (111) surfaces of a fcc crystal. *Left:* (110), (001) and (111) cuts through the crystal unit cell and atomic arrangement in the plane of the section. *Right:* Surface Brillouin zones corresponding to the cuts on the left-hand side [8.23]

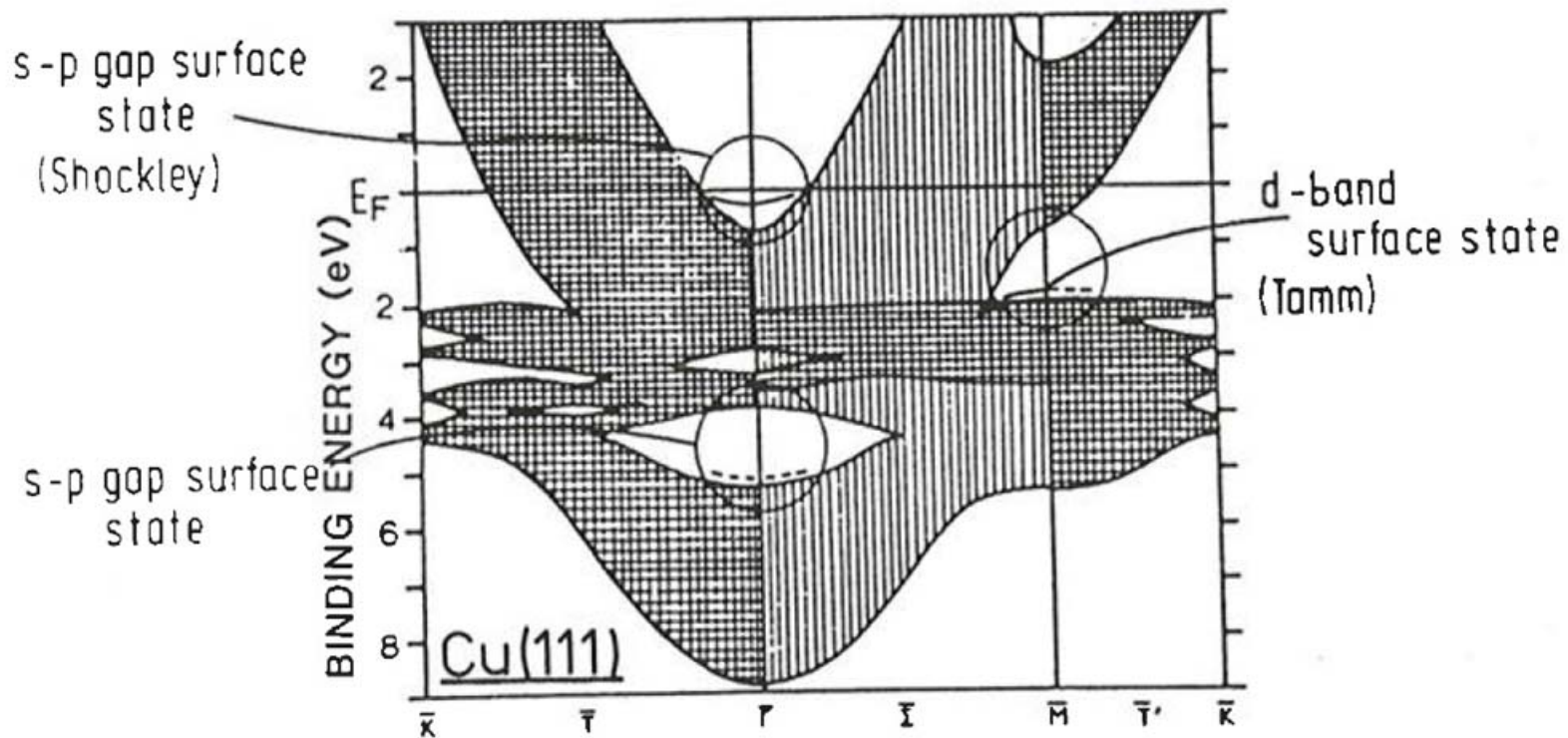


Fig.8.6. Density of states of Cu projected onto the (111) plane [8.23]. This provides the gaps in which a surface state can be expected to show up: near $\bar{\Gamma}$ with binding energies $E_F \leq E_s \leq 1\text{eV}$, near $\bar{\Gamma}$ with $3.7\text{eV} \leq E_s \leq 5.1\text{eV}$ and near \bar{M} with $1\text{eV} \leq E_s \leq 2\text{eV}$

