

Ring-Hamiltonian

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1 electron on a ring of protons at sites R_1, R_2, \dots, R_N

- equivalent intersite interval a

- length $L = N \cdot a$

- define quantum number $k \equiv \frac{2\pi}{L} \cdot n$ with $n = 1, \dots, N$
or $n = -\frac{N}{2} + 1, \dots, 0, \dots, \frac{N}{2}$

- define wavefunction:

$$\Psi_k(r) \equiv \frac{1}{\sqrt{N}} \sum_{R_i} e^{ikR_i} \varphi(r - R_i)$$

with

$$\langle \varphi(r - R_i) | H | \varphi(r - R_i) \rangle = E_0 \quad \text{or} \quad \langle \varphi(r) | H | \varphi(r) \rangle = E_0$$

$$\langle \varphi(r - R_i) | H | \varphi(r - R_{i+1}) \rangle = t \quad \text{or} \quad \langle \varphi(r) | H | \varphi(r+a) \rangle = t$$

$$\langle \varphi(r - R_i) | H | \varphi(r - R_{i-1}) \rangle = t \quad \text{or} \quad \langle \varphi(r) | H | \varphi(r-a) \rangle = t$$

and 0 otherwise

$$\begin{aligned} * \langle \Psi_k(r) | H | \Psi_k(r) \rangle &= \left\langle \frac{1}{\sqrt{N}} \sum_{R_i} e^{ikR_i} \varphi(r - R_i) \middle| H \middle| \frac{1}{\sqrt{N}} \sum_{R_j} e^{ikR_j} \varphi(r - R_j) \right\rangle \\ &= \frac{1}{N} \sum_{R_i} \sum_{R_j} e^{ik(R_j - R_i)} \langle \varphi(r - R_i) | H | \varphi(r - R_j) \rangle \end{aligned}$$

$$R \equiv R_j - R_i \quad = \frac{1}{N} \cdot N \cdot \sum_R e^{ikR} \langle \varphi(r) | H | \varphi(r - R) \rangle$$

$$\begin{aligned} R=0 &= 1 \cdot E_0 \\ R=a &= e^{ika} \cdot t \\ R=-a &= e^{-ika} \cdot t + \end{aligned}$$

$$= E_0 + t(e^{ika} + e^{-ika})$$

$$= E_0 + 2t \cos ka$$

$$= E_0 + 2t \cos\left(\frac{2\pi \cdot n}{N}\right)$$

$k \neq k'$:

$$\begin{aligned}
 \langle \psi_k(r) | H | \psi_{k'}(r) \rangle &= \left\langle \frac{1}{\sqrt{N}} \sum_{R_i} e^{ikR_i} \psi(r-R_i) \middle| H \middle| \frac{1}{\sqrt{N}} \sum_{R_j} e^{ik'R_j} \psi(r-R_j) \right\rangle \\
 &= \frac{1}{N} \sum_{R_i} \sum_{R_j} e^{-ikR_i} e^{ik'R_j} \langle \psi(r-R_i) | H | \psi(r-R_j) \rangle \\
 e^{ik'R_i} e^{-ik'R_i} = 1 &\rightarrow = \frac{1}{N} \sum_{R_i} \sum_{R_j} e^{i(k'-k)R_i} e^{ik'(R_j-R_i)} \langle \dots \dots \dots \rangle \\
 &= \frac{1}{N} \sum_{R_i} e^{i(k'-k)R_i} \sum_{R_j} e^{ik'(R_j-R_i)} \langle \dots \dots \dots \rangle \\
 * &= \frac{1}{N} \sum_{R_i} e^{i(k'-k)R_i} \sum_{R_j} e^{ik'(R_j-R_i)} \cdot f(R_j-R_i) \\
 ** &= \frac{1}{N} \sum_{R_i} e^{i(k'-k)R_i} \cdot F(k') \\
 &= \frac{F(k')}{N} \sum_{R_i} e^{i(k'-k) \cdot R_i}
 \end{aligned}$$

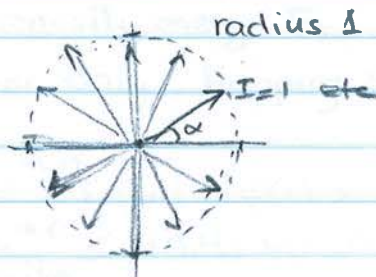
$k' = \frac{2\pi}{L}n', k = \frac{2\pi}{L}n, L = aN, m \equiv n' - n, R_i = a \cdot I$

$$= \frac{F(k')}{N} \sum_I e^{i \frac{2\pi}{N} \cdot m \cdot I}$$

$= 0$

→ wavefunctions with different k -values do not couple!

Angle $\alpha \equiv \frac{2\pi m}{N} \rightarrow$



The sum of all vectors $e^{i \frac{2\pi m}{N} \cdot I}$ in the complex plane is zero.

* $\langle \psi(r-R_i) | H | \psi(r-R_j) \rangle$ depends only on $(R_i - R_j)$, not on the absolute positions R_i or R_j

** $\sum_{R_j} e^{ik'(R_j-R_i)} \cdot f(R_j-R_i)$ does not depend on R_i due to periodicity.

Ring Hamiltonian in 1, 2, and 3 dimensions (or an infinitely large lattice in 1, 2, 3 dimensions)

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- assume only one basis function $\varphi(\underline{r})$ per site
- define Wannier function $\psi_{\underline{k}}(\underline{r}) = \frac{1}{\sqrt{N}} \sum_{\underline{R}} e^{i\underline{k} \cdot \underline{R}} \varphi(\underline{r} - \underline{R})$

\underline{R} covers all lattice sites

\underline{k} covers all points in the Brillouin Zone.

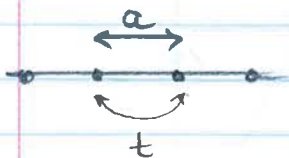
1 dim: $\underline{k} = k$ 2 dim: $\underline{k} = \begin{pmatrix} k_x \\ k_y \end{pmatrix}$ 3 dim: $\underline{k} = \begin{pmatrix} k_x \\ k_y \\ k_z \end{pmatrix}$

- for each \underline{k} one can calculate:

$$\begin{aligned} \langle \psi_{\underline{k}}(\underline{r}) | H | \psi_{\underline{k}}(\underline{r}) \rangle &= \left\langle \frac{1}{\sqrt{N}} \sum_{\underline{R}_1} e^{i\underline{k} \cdot \underline{R}_1} \varphi(\underline{r} - \underline{R}_1) \middle| H \middle| \frac{1}{\sqrt{N}} \sum_{\underline{R}_2} e^{i\underline{k} \cdot \underline{R}_2} \varphi(\underline{r} - \underline{R}_2) \right\rangle \\ &= \frac{1}{N} \sum_{\underline{R}_1} \sum_{\underline{R}_2} e^{i\underline{k} \cdot (\underline{R}_2 - \underline{R}_1)} \langle \varphi(\underline{r} - \underline{R}_1) | H | \varphi(\underline{r} - \underline{R}_2) \rangle \\ &= \sum_{\underline{R}} e^{i\underline{k} \cdot \underline{R}} \langle \varphi(\underline{r}) | H | \varphi(\underline{r} - \underline{R}) \rangle \end{aligned}$$

with \underline{R}_1 and \underline{R}_2 covering the same lattice sites
and with $\underline{R} \equiv \underline{R}_2 - \underline{R}_1$

Example 1: 1 dimensional, on-site energy E_0 ,
nearest neighbor only: hopping integral t

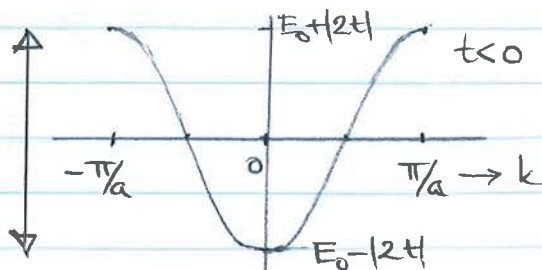


$$\begin{aligned} R=0 &\rightarrow e^{ik \cdot 0} \langle \varphi(\underline{r}) | H | \varphi(\underline{r}) \rangle = E_0 \\ R=a &\rightarrow e^{ika} \langle \varphi(\underline{r}) | H | \varphi(\underline{r}-a) \rangle = e^{ika} t \\ R=-a &\rightarrow e^{-ika} \langle \varphi(\underline{r}) | H | \varphi(\underline{r}+a) \rangle = e^{-ika} t \end{aligned}$$

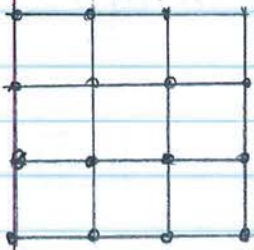
$$\langle \psi_{\underline{k}}(\underline{r}) | H | \psi_{\underline{k}}(\underline{r}) \rangle = E_0 + 2t \cos(ka)$$

$$\begin{aligned} ka &= \frac{2\pi}{N} \cdot n \\ &= \frac{2\pi}{a} \cdot \frac{n}{N} \end{aligned}$$

$$\text{bandwidth} = |4t|$$



example 2 : 2 dimensional square lattice
on-site energy E_0 , nearest neighbor hopping t



$$\underline{R} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} + a \begin{pmatrix} n_x \\ n_y \end{pmatrix} \quad n_x, n_y = \text{Integer}$$

$$\underline{k} = \begin{pmatrix} k_x \\ k_y \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} + \frac{2\pi}{a} \cdot \frac{1}{N} \begin{pmatrix} n_x \\ n_y \end{pmatrix}$$

$$\underline{R} = \underline{0} \rightarrow e^{i\mathbf{k}\cdot\underline{0}} \langle \varphi(\underline{r}) | H | \varphi(\underline{0}) \rangle = E_0$$

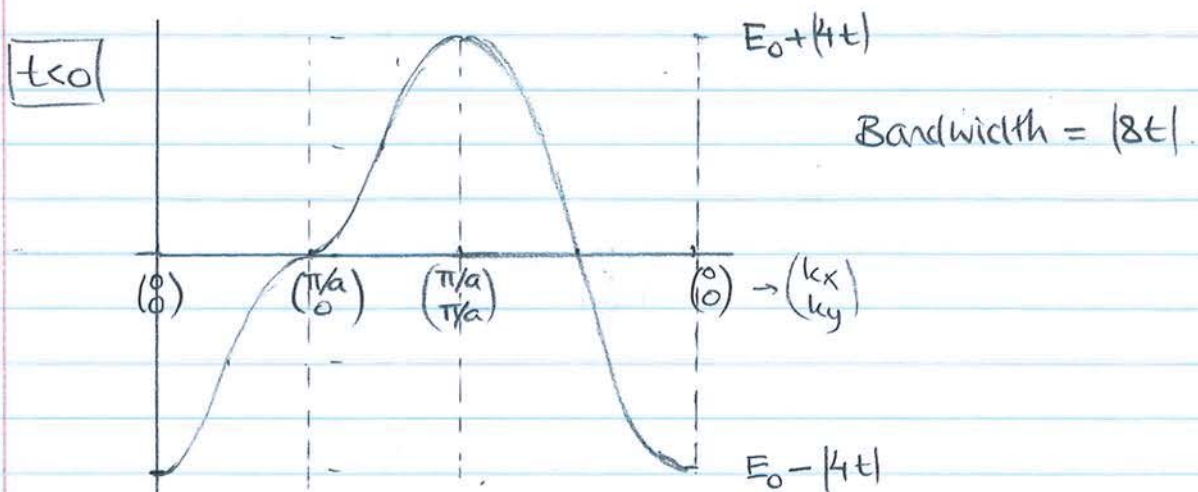
$$\underline{R} = \begin{pmatrix} a \\ 0 \end{pmatrix} \rightarrow e^{ik_x a} \langle \varphi(\underline{r}) | H | \varphi(\underline{r} - \begin{pmatrix} a \\ 0 \end{pmatrix}) \rangle = e^{ik_x a} \cdot t$$

$$\underline{R} = \begin{pmatrix} -a \\ 0 \end{pmatrix} \rightarrow e^{-ik_x a} \langle \varphi(\underline{r}) | H | \varphi(\underline{r} - \begin{pmatrix} -a \\ 0 \end{pmatrix}) \rangle = e^{-ik_x a} \cdot t$$

$$\underline{R} = \begin{pmatrix} 0 \\ a \end{pmatrix} \rightarrow e^{ik_y a} \langle \varphi(\underline{r}) | H | \varphi(\underline{r} - \begin{pmatrix} 0 \\ a \end{pmatrix}) \rangle = e^{ik_y a} \cdot t$$

$$\underline{R} = \begin{pmatrix} 0 \\ -a \end{pmatrix} \rightarrow e^{-ik_y a} \langle \varphi(\underline{r}) | H | \varphi(\underline{r} - \begin{pmatrix} 0 \\ -a \end{pmatrix}) \rangle = e^{-ik_y a} \cdot t$$

$$\langle \psi_{\underline{k}}(\underline{r}) | H | \psi_{\underline{k}}(\underline{r}) \rangle = E_0 + 2t \{ \cos(k_x a) + \cos(k_y a) \}$$



example 3: 3 dimensional simple cubic lattice
 on-site energy E_0 , nearest neighbor hopping t

$$\langle \psi_{\mathbf{k}} | H | \psi_{\mathbf{k}} \rangle = E_0 + 2t \{ \cos(k_x a) + \cos(k_y a) + \cos(k_z a) \}$$

